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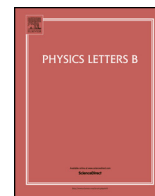
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Landau pole in the Standard Model with weakly interacting scalar fields



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ABSTRACT

We consider the Standard Model with a new scalar field X which is an n_X representation of the $SU(2)_L$ with a hypercharge Y_X . The renormalization group running effects on the new scalar quartic coupling constants are evaluated. Even if we set the scalar quartic coupling constants to be zero at the scale of the new scalar field, the coupling constants are induced by the one-loop effect of the weak gauge bosons. Once non-vanishing couplings are generated, the couplings rapidly increase by renormalization group effect of the quartic coupling constant itself. As a result, the Landau pole appears below Planck scale if $n_X \geq 4$. We find that the scale of the obtained Landau pole is much lower than that evaluated by solving the one-loop beta function of the gauge coupling constants.

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The discovery of the Higgs boson [1,2] and the determination of its mass open up the possibility that the Standard Model (SM) can be valid up to very high energy scale such as string/Planck scale [3–16].¹ In particular, the Higgs self-coupling constant and its beta function almost vanish at the same time around Planck scale if the top-quark mass M_t is 171 GeV. This fact may imply relations between physics at the weak scale and physics at Planck scale, e.g., Higgs inflation [26–33], a multiple point criticality principle [34] and a maximum entropy principle [35–39].

On the other hand, we know that the SM should be extended to include dark matter, the origin of neutrino masses, and a mechanism for generating the baryon number of the universe. Such an extension often introduces a new scalar field X at around the electroweak/TeV or some intermediate scale which is charged under the $SU(2)_L \times U(1)_Y$ gauge group. If we add a sufficiently large $SU(2)_L$ isospin multiplet, the new field provides a good dark matter candidate because direct interactions with SM particles are forbidden automatically by renormalizability [40–42]. It is known as the Type-II seesaw mechanism that Majorana masses for neutrinos can be generated by introducing a vacuum expectation value of a scalar triplet with $Y = 1$ [43–46]. It is also known that additional scalar boson loops can increase the triple (SM-like) Higgs boson coupling [47,48], which helps to satisfy the sphaleron decoupling condition for the electroweak baryogenesis [49–51]. A special choice of the Higgs multiplet, the Higgs septet with $Y = 2$ [52,53], can give

sizable deviations in the observed Higgs boson couplings without conflicting the stringent constraint from the electroweak ρ parameter. There are so many extended models motivated by various motivations, so that it is important to constrain such possibilities in a generic way.

The triviality bound [54,55] is studied to give such a generic constraint on the extended Higgs models [56,57]. Considering the renormalization group equation (RGE) for the Higgs quartic coupling constant, the energy dependent coupling constant $\lambda(Q^2)$ grows with an energy scale Q . At the end of the day, the running coupling constant blows up at a certain scale Λ_{LP} . The scale is called Landau pole (LP), where the coupling constant becomes infinite $1/\lambda(\Lambda_{LP}) = 0$. The absence of the LP at a certain scale, e.g., Planck scale, is often required to constrain the Higgs quartic coupling (equivalently the Higgs boson mass in the SM) [54,55]. This bound can be understood as a criterion for perturbativity of the theory.

In this letter, we give the RGE analysis of the scalar quartic coupling constants in the SM with one more scalar multiplet X , where the field X is an n_X representation under $SU(2)_L$ with a hypercharge Y_X . The model predicts the LP below Planck scale for $n_X \geq 4$ if X appears in the electroweak/TeV scale, where the scale of the LP is defined by the blowup of a coupling constant. The LP we will derive from the scalar quartic coupling constants is much lower than that obtained by solving one-loop beta functions of the gauge coupling constants [58]. The point is that the quartic coupling constants of X are rapidly induced by the electroweak gauge couplings with large coefficients (due to the large electroweak charges) in the beta functions, even if the initial values

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¹ See Refs. [17–25] for the analysis in simple extensions of the SM.

Table 1

The independent coupling constants in the scalar potential are listed. We also show the dimension four extra couplings which contribute to the beta functions.

(n_X, Y_X)	independent couplings	dim-4 extra couplings
(4, 1/2)	$\lambda, \kappa, \kappa', \lambda_X, \lambda'_X$	$\Phi^{\dagger 2} X^2, \Phi \Phi^{\dagger 2} X$
(4, 3/2)	$\lambda, \kappa, \kappa', \lambda_X, \lambda'_X$	$\Phi^{\dagger 3} X$
(5, real)	$\lambda, \kappa, \lambda_X$	\times
(5, 1)	$\lambda, \kappa, \kappa', \lambda_X, \lambda'_X, \lambda''_X$	\times
(5, 2)	$\lambda, \kappa, \kappa', \lambda_X, \lambda'_X, \lambda''_X$	\times
(6, 1/2)	$\lambda, \kappa, \kappa', \lambda_X, \lambda'_X, \lambda''_X$	$\Phi^{\dagger 2} X^2$
(6, 3/2)	$\lambda, \kappa, \kappa', \lambda_X, \lambda'_X, \lambda''_X$	\times
(6, 5/2)	$\lambda, \kappa, \kappa', \lambda_X, \lambda'_X, \lambda''_X$	\times
(7, real)	$\lambda, \kappa, \lambda_X, \lambda''_X$	\times
(7, 1)	$\lambda, \kappa, \kappa', \lambda_X, \lambda'_X, \lambda''_X, \lambda'''_X$	\times
(7, 2)	$\lambda, \kappa, \kappa', \lambda_X, \lambda'_X, \lambda''_X, \lambda'''_X$	\times
(7, 3)	$\lambda, \kappa, \kappa', \lambda_X, \lambda'_X, \lambda''_X, \lambda'''_X$	\times

of the quartic coupling constants are set to be zeros. Once a finite quartic coupling is injected, the RGE running of the quartic coupling constant leads the LP. In general, there are many degrees of freedom to choose the initial condition for the new coupling constants. Among them, we evaluate a conservative scale (the largest scale) of the LP.

The scalar potential of the model is given by

$$V = -\mu^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2 + M_X^2 X^\dagger X + \lambda_X |X^\dagger X|^2 + \kappa |X^\dagger X| |\Phi^\dagger \Phi| + \kappa' (X^\dagger T_X^a X) (\Phi^\dagger T_\Phi^a \Phi) + \lambda'_X (X^\dagger T_X^a X)^2 + \lambda''_X (X^\dagger T_X^a T_X^b X)^2 + \dots, \quad (1)$$

where Φ is the SM Higgs doublet, X is a new scalar field, and the $SU(2)_L$ generator for X is T_X^a ($a = 1, 2, 3$). Among the scalar quartic coupling constants, some of them are related to each other depending on the electroweak charges of X . For quadruplets, extra coupling constants are allowed,

$$\lambda_{\Phi \Phi^{\dagger 2} X} \Phi \Phi^\dagger \Phi^\dagger X = \lambda_{\Phi \Phi^{\dagger 2} X} \Phi^i \Phi_j^\dagger \Phi_k^\dagger X^{ijk} \epsilon_{ii'}, \quad (2)$$

$$\lambda_{\Phi^{\dagger 2} X^2} \Phi^{\dagger 2} X^2 = \lambda_{\Phi^{\dagger 2} X^2} \Phi_i^\dagger \Phi_j^\dagger X^{ik\ell} X^{jk\ell'} \epsilon_{kk'} \epsilon_{\ell\ell'}, \quad (3)$$

$$\lambda_{\Phi^{\dagger 3} X} \Phi^{\dagger 3} X = \lambda_{\Phi^{\dagger 3} X} \Phi_i^\dagger \Phi_j^\dagger \Phi_k^\dagger X^{ijk}, \quad (4)$$

where we adopt the symmetric tensor notation, i.e., $(X^{111}, X^{112}, X^{122}, X^{222}) = (X^1, X^2/\sqrt{3}, X^3/\sqrt{3}, X^4)$ for $n_X = 4$. The former two coupling constants exist for $Y_X = 1/2$, while the last does for $Y_X = 3/2$. There is a similar coupling of $\lambda_{\Phi^{\dagger 2} X^2}$ for a sextet. Note that there is an accidental global $U(1)$ symmetry, if all the additional dimension four couplings are forbidden. We summarize the independent coupling constants and possible dimension four extra coupling constants of each model in Table 1.

Next, let us move to the RGE analysis. The scale below the mass of X , M_X , we use the SM beta functions. Above M_X , the runnings of the electroweak gauge couplings are modified to as

$$\frac{dg_Y}{dt} = \frac{1}{16\pi^2} \left(\frac{41}{6} + \frac{1}{6} A Y_X^2 n_X \right) g_Y^3, \quad (5)$$

$$\frac{dg_2}{dt} = \frac{1}{16\pi^2} \left(-\frac{19}{6} + \frac{1}{6} A T(X) \right) g_2^3, \quad (6)$$

where $t = \ln \mu$ with μ being the renormalization scale, $A = 1$ (2) for a real (complex) scalar field, $T(X)$ is the Dynkin index. We see that g_2 is non-asymptotic free if $n_X \geq 7$ (5) for a real (complex) field case. The runnings of the top-Yukawa coupling y_t and the strong gauge coupling g_3 are unchanged at one-loop level. The one-loop RGEs of the scalar coupling constants are given in Appendix A.

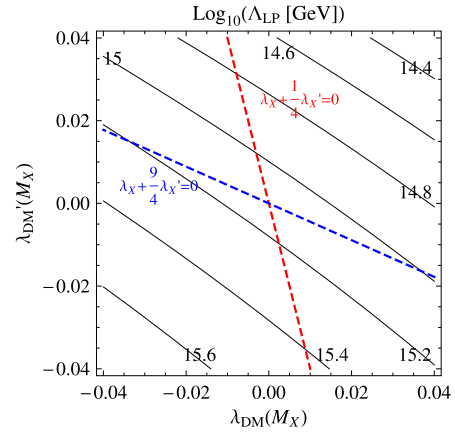


Fig. 1. A contour plot of the scale of the LP as functions of λ_X, λ'_X at $\mu = M_X$ with $M_X = 100$ GeV. The red and blue lines represent to $\lambda_X + 1/4 \lambda'_X = 0$ and $\lambda_X + 9/4 \lambda'_X = 0$, respectively. The left-bottom side of the red or blue line corresponds to the unbounded scalar potential.

As a set of initial conditions for the RGEs, we take

$$\lambda_X(M_X) = \lambda'_X(M_X) = \dots = \kappa(M_X) = \kappa'(M_X) = 0, \quad (7)$$

and evaluate the scale of the LP. This set gives a conservative evaluation for the scale of the LP under the condition that the scalar potential is bounded from below. We use the initial values of SM parameters in Ref. [12] and take $M_H = 125.7$ GeV, $M_t = 173$ GeV, which are not sensitive to our analysis. By calculating the RGE, the existence of the LP is found below Planck scale for $n_X \geq 4$. The blowup of the quartic coupling constants $\lambda_X, \lambda'_X, \dots$ of X occurs because the quantum corrections to these couplings are proportional to the fourth power of the electroweak charges of X . This gives a rapid injection of non-zero value to the quartic coupling constants, and eventually results in the LP.

Let us discuss the initial condition for the new quartic couplings constants. For an illustration, we focus on the case $(n_X, Y_X) = (4, 3/2)$ with an additional Z_2 symmetry, which forbids the term $\Phi^{\dagger 3} X$. In this case, by imposing the boundedness of the scalar potential for $|X^1|^4$ and $|X^2|^4$ directions, we obtain the following conditions²

$$\lambda_X + \frac{1}{4} \lambda'_X \geq 0, \quad \lambda_X + \frac{9}{4} \lambda'_X \geq 0. \quad (8)$$

This requirements comes from the $|X^1|^4$ and $|X^2|^4$ terms in the scalar potential of X ,

$$\begin{aligned} \lambda_X |X^\dagger X|^2 + \lambda'_X (X^\dagger T^a X)^2 \\ = \lambda_X (|X^1|^4 + |X^2|^4) + \lambda'_X \left(\frac{9}{4} |X^1|^4 + \frac{1}{4} |X^2|^4 \right) + \dots, \end{aligned} \quad (9)$$

where the second term in the r.h.s. of the equation results from $(X^\dagger T^3 X)^2$ with $T^3 = \text{diag}(3/2, 1/2, -1/2, -3/2)$. The same conditions are also come from $|X^3|^4$ and $|X^4|^4$ directions. The contour plot for the scale of the LPs together with the above two conditions is shown in Fig. 1 as functions of $\lambda_X(M_X)$ and $\lambda'_X(M_X)$. One can see that $\lambda_X = \lambda'_X = 0$ actually gives the conservative scale (the largest scale) of the LP. Even if we introduce nonzero $\kappa, \kappa', \lambda_{\Phi^{\dagger 3} X}$, these effects decrease the scale of the LP since these contributions to the beta functions of λ_X and λ'_X are always positive. For the $(n_X, Y_X) = (4, 1/2)$ case, the similar arguments are applicable for

² Note that this is not the sufficient condition for the boundedness but the necessary condition.

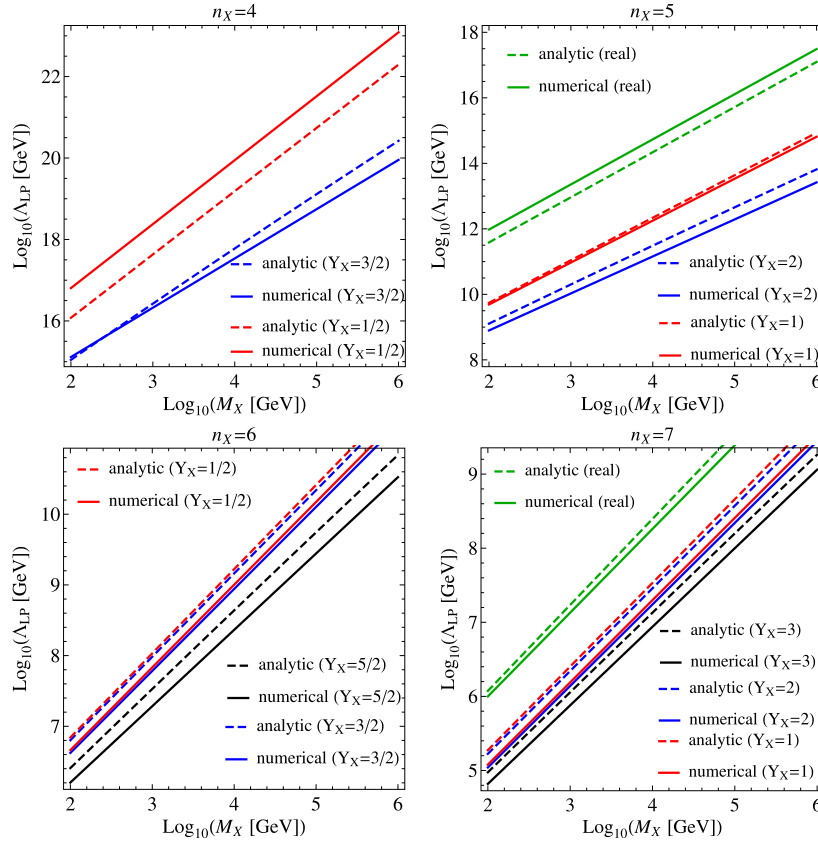


Fig. 2. The solid and dashed lines correspond to the numerical results and the approximated results in Eq. (14).

λ_X, λ'_X and $\lambda_{\Phi\Phi^{\dagger 2}X}$. Although the effect of $\lambda_{\Phi\Phi^{\dagger 2}X}$ is unclear, we numerically confirmed that this coupling constant also leads to the lower scale of the LP. Thus, we conclude that Eqs. (7) are the initial condition for the conservative evaluation of the LP. The same discussions would be applied to higher isospin multiplets. We also confirmed these points numerically. The detailed and generalized analysis will be presented in our future publications [59].

In Fig. 2, we present the conservative scale of the LP as a function of M_X in the SM with one more scalar multiplet ($n_X = 4, 5, 6, 7$).³ We consider all the possible assignments of the hypercharges, where fractional electric charges are not allowed for new particles. For a $Y_X = 0$ field, a real scalar condition is assumed. The solid lines in each plot represent the numerical results, while the dashed lines show the approximated results obtained from the analytic study as we will discuss later in this letter. We fit the solid lines in Fig. 2 by the scale of the LP and M_X with an exponent. The results are listed in Table 2. For comparisons, we also present the positions of the LP Λ_{gi} which are calculated by solving the one-loop beta function of gauge couplings.

$$\begin{aligned}\Lambda_{gi} &= M_X \exp\left(\frac{1}{2B_i} \frac{1}{g_i^2(M_X)}\right) \\ &= M_X \left(\frac{M_t}{M_X}\right)^{B_{i,\text{SM}}/B_i} \exp\left(\frac{1}{2B_i} \frac{1}{g_i^2(M_t)}\right),\end{aligned}\quad (10)$$

where $i = Y, 2$, and $B_{i,\text{SM}}$ is the coefficient of the beta function of g_i including a factor $1/16\pi^2$. We see that the conservative scale

³ Requiring the tree-level unitarity of $SU(2)_L$ gauge interactions, $n_X \leq 8(9)$ is obtained for a complex (real) scalar multiplet [60,61].

Table 2

The fitted results of the scale of LPs with the initial condition in Eqs (7). The positions of LP derived from one-loop beta function of the gauge couplings taking $M_X = 100$ GeV are also given.

(n_X, Y_X)	LP [GeV]	Λ_{g_Y} [GeV]	Λ_{g_2} [GeV]
(4, 1/2)	$6.4 \times 10^{16} (M_X/100 \text{ GeV})^{1.57}$	1.2×10^{41}	–
(4, 3/2)	$1.3 \times 10^{15} (M_X/100 \text{ GeV})^{1.21}$	3.1×10^{30}	–
(5, real)	$9.5 \times 10^{11} (M_X/100 \text{ GeV})^{1.38}$	9.6×10^{42}	–
(5, 1)	$4.9 \times 10^9 (M_X/100 \text{ GeV})^{1.28}$	9.0×10^{34}	9.0×10^{488}
(5, 2)	$7.9 \times 10^8 (M_X/100 \text{ GeV})^{1.13}$	5.7×10^{22}	9.0×10^{488}
(6, 1/2)	$4.6 \times 10^6 (M_X/100 \text{ GeV})^{1.17}$	1.6×10^{40}	2.7×10^{32}
(6, 3/2)	$4.2 \times 10^6 (M_X/100 \text{ GeV})^{1.16}$	5.2×10^{26}	2.7×10^{32}
(6, 5/2)	$1.6 \times 10^6 (M_X/100 \text{ GeV})^{1.08}$	3.2×10^{16}	2.7×10^{32}
(7, real)	$1.0 \times 10^6 (M_X/100 \text{ GeV})^{1.13}$	9.6×10^{42}	1.3×10^{56}
(7, 1)	$1.2 \times 10^5 (M_X/100 \text{ GeV})^{1.11}$	3.6×10^{32}	1.4×10^{15}
(7, 2)	$1.1 \times 10^5 (M_X/100 \text{ GeV})^{1.10}$	2.2×10^{19}	1.4×10^{15}
(7, 3)	$6.6 \times 10^4 (M_X/100 \text{ GeV})^{1.06}$	1.2×10^{12}	1.4×10^{15}

of the LP from the scalar quartic coupling constant is much lower than Λ_{gi} .

In order to understand the results analytically, we focus on the beta function of the quartic coupling of X in which the gauge coupling $g(\mu)$ is approximated by the constant $g(M_X)$. Then, we redefine the quartic coupling constants λ'_X, \dots such that g^4 terms only appear in one beta function. By this redefinition, we correctly take into account the induced quartic coupling constant from the gauge couplings at one-loop level, since we begin with a conservative assumption given in Eqs. (7). We call the new coupling constants $\tilde{\lambda}'_X, \dots$. Consequently, the RGE is simplified as

$$\frac{d\lambda_X}{dt} = \frac{1}{16\pi^2} (c_0 - c_1\lambda_X + c_2\lambda_X^2 + c_3\lambda_X\tilde{\lambda}'_X + c_4\tilde{\lambda}'_X{}^2), \quad (11)$$

$$\frac{d\tilde{\lambda}'_X}{dt} = \frac{1}{16\pi^2} (-\tilde{c}'_1\lambda_X + \tilde{c}'_2\lambda_X^2 + \tilde{c}'_3\lambda_X\tilde{\lambda}'_X + \tilde{c}'_4\tilde{\lambda}'_X{}^2). \quad (12)$$

However, even if we set $\tilde{\lambda}'_X = 0$ at M_X , nonzero $\tilde{\lambda}'_X$ is induced by the \tilde{c}'_2 term, which affects the running of λ_X through the c_3 term.⁴ Therefore, we redefine λ_X in order to cancel the c_3 term. As a result, we have

$$\frac{d\tilde{\lambda}_X}{dt} = \frac{1}{16\pi^2} (\tilde{c}_0 - \tilde{c}_1 \tilde{\lambda}_X + \tilde{c}_2 \tilde{\lambda}_X^2), \quad (13)$$

taking $\tilde{\lambda}'_X = \dots = 0$. This simplified differential equation can be solved analytically. The position of the LP is calculated as

$$\frac{\Lambda_{\text{LP}}}{M_X} = \exp \left[\frac{16\pi^2}{\tilde{c}_2 d} \left\{ \frac{\pi}{2} - \tan^{-1} \left[\frac{1}{d} \left(\tilde{\lambda}_X(M_X) - \frac{\tilde{c}_1}{2\tilde{c}_2} \right) \right] \right\} \right], \quad (14)$$

$$d = \sqrt{\frac{\tilde{c}_0}{\tilde{c}_2} - \frac{\tilde{c}_1^2}{4\tilde{c}_2^2}}. \quad (15)$$

As an example, we show the concrete procedure for the case with $n_X = 4$. Concentrating on the quartic coupling constants of X , we have

$$\frac{d\lambda_X}{dt} = \frac{1}{16\pi^2} \left(+32\lambda_X^2 + 30\lambda_X\lambda'_X + \frac{99}{2}\lambda'^2_X + 6Y_X^4 g_Y^4 + \frac{297}{8}g_2^4 - 12Y_X^2 \lambda_X g_Y^2 - 45\lambda_X g_2^2 \right), \quad (16)$$

$$\frac{d\lambda'_X}{dt} = \frac{1}{16\pi^2} \left(+24\lambda_X\lambda'_X + 42\lambda'^2_X + 12Y_X^2 g_Y^2 g_2^2 - 12Y_X^2 \lambda'_X g_Y^2 - 45\lambda'_X g_2^2 \right). \quad (17)$$

Here we put $g_Y(\mu) = g_Y(M_X)$, $g_2(\mu) = g_2(M_X)$ and $\kappa = \kappa' = \lambda_{\Phi\Phi^\dagger X} = \lambda_{\Phi^2 X^2} = \lambda_{\Phi^3 X} = \lambda_{\Phi^4 X} = 0$. At first, we redefine λ'_X in order to cancel the constant term ($12Y_X^2 g_Y^2 g_2^2$ term). Hence, we take

$$\tilde{\lambda}'_X = \lambda'_X + \eta \lambda_X, \quad (18)$$

where $\eta \equiv -12Y_X^2 g_Y^2 g_2^2 / (6Y_X^4 g_Y^4 + \frac{297}{8}g_2^4)$. The requirement of the vanishing $g_Y^2 g_2^2$ term fixes $\tilde{\lambda}'_X$ up to the normalization, while we still have the degree of freedom to redefine λ_X . Using this, we define $\tilde{\lambda}_X$ as

$$\tilde{\lambda}_X = \lambda_X + \xi \tilde{\lambda}'_X, \quad (19)$$

and choose ξ in such a way that the $\tilde{\lambda}_X \tilde{\lambda}'_X$ term vanishes in the beta function of $\tilde{\lambda}_X$. Then, we obtain

$$\xi = \frac{-20 + 3\eta - 99\eta^2}{16\eta + 24\eta^2 + 99\eta^3} \pm \frac{\sqrt{5(80 + 72\eta + 621\eta^2)}}{16\eta + 24\eta^2 + 99\eta^3}. \quad (20)$$

The coefficients c_0, c_1, c_2 are given by

$$\tilde{c}_0 = \frac{297}{8}g_2^4 + 6Y_X^4 g_Y^4, \quad (21)$$

$$\tilde{c}_1 = 12Y_X^2 g_Y^2 + 45g_2^2, \quad (22)$$

$$\tilde{c}_2 = 32 - 30\eta + \frac{99}{2}\eta^2 + 8\eta\xi + 12\eta^2\xi + \frac{99}{2}\eta^3\xi. \quad (23)$$

By this procedure, we can evaluate the LP in the analytical calculations. The approximated lines are also plotted in Fig. 2. Comparing the numerical results and our analytic results in these figures, the above expressions fit the numerical results within an order of magnitude error. From this analytical analysis, we confirm that the

finite quartic coupling constants are originally generated by the gauge couplings, and then the RGE of the quartic coupling constant leads the LP. Roughly speaking, the LP is induced by two-loop radiative corrections to the RGEs. We can also evaluate the exponent of M_X by putting $g_{Y(2)}(M_X) \simeq g_{Y(2)}(M_t) + B_{Y(2),\text{SM}} \ln(M_X/M_t)$. Although the analytic expression is too complicate to write, we obtain ~ 1.52 and ~ 1.28 for $(4, 1/2)$ and $(4, 3/2)$, respectively.

Before concluding this letter, a few more remarks are given in order. Firstly, the position of the LP does not depend on the normalization of $\tilde{\lambda}_X$. We can understand the reason from Eqs. (13) and (14) as follows; let us define the new scalar coupling constant having a different normalization as $\tilde{\lambda}_X^N := a\tilde{\lambda}_X$ where a is a real constant. Although the RGE of $\tilde{\lambda}_X^N$ takes the same form as $\tilde{\lambda}_X$, the coefficients are different. They are ac_0, c_1 and c_2/a . By using these coefficients and $\tilde{\lambda}_X^N(M_X) = a\tilde{\lambda}_X(M_X)$, one can easily show that the LP for $\tilde{\lambda}_X^N$ calculated by Eq. (14) gives the same result as that of $\tilde{\lambda}_X$. Therefore, we do not need to care about the normalizations of scalar couplings when we study the LP.

Secondly, we discuss the nonzero value of the initial condition. Let us start with Eq. (14). Expanding Eq. (14) around $1/\tilde{\lambda}_X(M_X) = 0$, we get

$$\frac{\Lambda_{\text{LP}}}{M_X} = \exp \left[\frac{16\pi^2}{\tilde{c}_2} \frac{1}{\tilde{\lambda}_X(M_X)} + \mathcal{O}(1/\tilde{\lambda}_X(M_X)^2) \right], \quad (24)$$

which means that the scale of the LP is determined by only the initial value and \tilde{c}_2 . The conditions to neglect higher order terms are $\tilde{\lambda}_X(M_X) \gtrsim \tilde{c}_1/(2\tilde{c}_2)$, $d, \tilde{c}_1/(2\tilde{c}_2) \approx 0.28, d \approx 0.31$ for $(n_X, Y_X) = (4, 3/2)$.

Though we have focused on a single scalar extension of the SM in this letter, more matter fields may be added. Since the increase of scalar fields leads to the large coefficient in the beta functions in general, we expect that the scale of the LP becomes lower. If we add fermions in addition to scalars, the Yukawa couplings among them may be allowed. In such a case, the scale of the LP can be large since the Yukawa coupling gives the negative contribution to the beta function at the one-loop level.

Finally, we comment on the implications of the new physics beyond the SM. In the Minimal dark matter models [40–42], interactions of new scalar multiplets are also bounded by the relic abundance of dark matter and the data for direct/indirect searches. It is very interesting to combine our study with such constraints [59]. In some class of the seesaw models for the neutrino mass generation, Majorana fermions are introduced together with scalar multiplets, which may or may not accommodate a dark matter candidate. These models are constrained not only by the relic abundance but also by the neutrino oscillation data. In these models, the effects of Majorana fermions on the RGEs are expected to be small, because the required Yukawa coupling is generally small if we add these new fermions in the electroweak/TeV scale. In addition, models with a flavor symmetry are also very intriguing candidates for this kind of analyses. Because the scalar fields are embedded in the multiplet under the flavor symmetry, which provides a large number of $SU(2)_L$ multiplets. Increase of the number of multiplet also gives large contribution to the RGE, which would make the LP smaller [59].

In conclusion, we have investigated the scale of the LP using the one-loop RGE in the SM with one more scalar n_X -plet X ($n_X = 4, 5, 6, 7$ with all possible hypercharge assignments). The LP is found below Planck scale for $n_X \geq 4$, if we introduce a new scalar multiplet at electroweak/TeV scale. This means that any single scalar field extension with $n_X \geq 4$ of the SM is not allowed, if we impose the absence of the LP up to Planck scale. The scale is evaluated with the conservative initial condition, where the initial values for the quartic coupling constant of new scalar fields are set

⁴ We checked that the effect of the c_4 term is smaller than that of the c_3 term.

to be zeros at the scale of new particles. Nevertheless, the quartic coupling constants are induced from the electroweak gauge interactions. The induced coupling constants are rapidly enhanced, and finally hit the LP. We have calculated the LP as a function of the mass of X , and listed a fitting formula of the LP with an exponent. The results are consistent with our approximated formula within an order of magnitude error, which is derived from the simplified RGE with the conservative assumption. The obtained LP in each model is much smaller than that previously calculated by solving the beta functions of the gauge couplings [58]. These new results are very useful and generic constraints on the beyond the SM including new scalar multiplets.

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Appendix A. Renormalization group equations

We calculate the RGEs by using the general formula in Ref. [62].

- SM with a quadruplet scalar field

$$\frac{d\lambda}{dt} = \frac{1}{16\pi^2} \left(+24\lambda^2 - 6y_t^4 + \frac{3}{8}g_Y^4 + \frac{9}{8}g_2^4 + \frac{3}{4}g_Y^2g_2^2 + 12\lambda y_t^2 - 3\lambda g_Y^2 - 9\lambda g_2^2 + 4\kappa^2 + \frac{5}{4}\kappa'^2 + \frac{40}{9}|\lambda_{\Phi^{\dagger 2}X^2}|^2 + 8|\lambda_{\Phi\Phi^{\dagger 2}X}|^2 + 24\lambda_{\Phi^{\dagger 3}X}^2 \right),$$

$$\frac{d\lambda_X}{dt} = \frac{1}{16\pi^2} \left(+32\lambda_X^2 + 30\lambda_X\lambda'_X + \frac{99}{2}\lambda_X'^2 + 6Y_X^4g_Y^4 + \frac{297}{8}g_2^4 - 12Y_X^2\lambda_Xg_Y^2 - 45\lambda_Xg_2^2 + 2\kappa^2 + 2|\lambda_{\Phi^{\dagger 2}X^2}|^2 \right),$$

$$\frac{d\lambda'_X}{dt} = \frac{1}{16\pi^2} \left(+24\lambda_X\lambda'_X + 42\lambda_X'^2 + 12Y_X^2g_Y^2g_2^2 - 12Y_X^2\lambda'_Xg_Y^2 - 45\lambda'_Xg_2^2 + \frac{1}{2}\kappa'^2 - \frac{8}{9}|\lambda_{\Phi^{\dagger 2}X^2}|^2 \right),$$

$$\frac{d\kappa}{dt} = \frac{1}{16\pi^2} \left(+3Y_X^2g_Y^4 + \frac{45}{4}g_2^4 + 4\kappa^2 + \frac{15}{4}\kappa'^2 + 6y_t^2\kappa - 6Y_X^2\kappa g_Y^2 - \frac{3}{2}\kappa g_Y^2 - 27\kappa g_2^2 + 12\kappa\lambda + 20\kappa\lambda_X + 15\kappa\lambda'_X + 18\lambda_{\Phi^{\dagger 3}X}^2 + \frac{40}{3}|\lambda_{\Phi^{\dagger 2}X^2}|^2 + 6|\lambda_{\Phi\Phi^{\dagger 2}X}|^2 \right),$$

$$\frac{d\kappa'}{dt} = \frac{1}{16\pi^2} \left(+12Y_Xg_Y^2g_2^2 + 6y_t^2\kappa' - \frac{3}{2}\kappa'g_Y^2 - 6Y_X^2\kappa'g_Y^2 - 27\kappa'g_2^2 + 8\kappa\kappa' + 4\kappa'\lambda + 4\kappa'\lambda_X + 31\kappa'\lambda'_X + 24\lambda_{\Phi^{\dagger 3}X}^2 + \frac{64}{9}|\lambda_{\Phi^{\dagger 2}X^2}|^2 + \frac{8}{3}|\lambda_{\Phi\Phi^{\dagger 2}X}|^2 \right),$$

$$\frac{d\lambda_{\Phi^{\dagger 2}X^2}}{dt} = \frac{1}{16\pi^2} \left(-4\lambda_{\Phi^{\dagger 2}X}^2 + 6y_t^2\lambda_{\Phi^{\dagger 2}X^2} - \frac{3}{2}\lambda_{\Phi^{\dagger 2}X^2}g_Y^2 - 6Y_X^2\lambda_{\Phi^{\dagger 2}X^2}g_Y^2 - 27\lambda_{\Phi^{\dagger 2}X^2}g_2^2 + 8\kappa\lambda_{\Phi^{\dagger 2}X^2} + 4\kappa'\lambda_{\Phi^{\dagger 2}X^2} + 4\lambda\lambda_{\Phi^{\dagger 2}X^2} + 4\lambda_X\lambda_{\Phi^{\dagger 2}X^2} - 11\lambda'_X\lambda_{\Phi^{\dagger 2}X^2} \right),$$

$$\frac{d\lambda_{\Phi\Phi^{\dagger 2}X}}{dt} = \frac{1}{16\pi^2} \left(+12\lambda\lambda_{\Phi\Phi^{\dagger 2}X} + 9y_t^2\lambda_{\Phi\Phi^{\dagger 2}X} - \frac{9}{4}g_Y^2\lambda_{\Phi\Phi^{\dagger 2}X} - 3Y_X^2g_Y^2\lambda_{\Phi\Phi^{\dagger 2}X} - 18g_2^2\lambda_{\Phi\Phi^{\dagger 2}X} + 6\kappa\lambda_{\Phi\Phi^{\dagger 2}X} + \frac{5}{2}\kappa'\lambda_{\Phi\Phi^{\dagger 2}X} - \frac{40}{3}\lambda_{\Phi^{\dagger 2}X^2}\lambda_{\Phi\Phi^{\dagger 2}X}^* \right),$$

$$\frac{d\lambda_{\Phi^{\dagger 3}X}}{dt} = \frac{\lambda_{\Phi^{\dagger 3}X}}{16\pi^2} \left(+12\lambda + 9y_t^2 - 18g_2^2 - \frac{9}{4}g_Y^2 - 3Y_X^2g_Y^2 + 6\kappa + \frac{15}{2}\kappa' \right). \quad (A.1)$$

- SM with a real quintet scalar field

$$\frac{d\lambda}{dt} = \frac{1}{16\pi^2} \left(+24\lambda^2 - 6y_t^4 + \frac{3}{8}g_Y^4 + \frac{9}{8}g_2^4 + \frac{3}{4}g_Y^2g_2^2 + \frac{5}{2}\kappa^2 + 12\lambda y_t^2 - 3\lambda g_Y^2 - 9\lambda g_2^2 \right),$$

$$\frac{d\lambda_X}{dt} = \frac{1}{16\pi^2} \left(+26\lambda_X^2 + 108g_2^4 - 72\lambda_Xg_2^2 + 2\kappa^2 \right),$$

$$\frac{d\kappa}{dt} = \frac{1}{16\pi^2} \left(+18g_2^4 + 12\kappa\lambda + 14\kappa\lambda_X + 6y_t^2\kappa - \frac{3}{2}\kappa g_Y^2 - \frac{81}{2}\kappa g_2^2 + 4\kappa^2 \right). \quad (A.2)$$

- SM with a complex quintet scalar field

$$\frac{d\lambda}{dt} = \frac{1}{16\pi^2} \left(+24\lambda^2 - 6y_t^4 + \frac{3}{8}g_Y^4 + \frac{9}{8}g_2^4 + \frac{3}{4}g_Y^2g_2^2 + 12\lambda y_t^2 - 3\lambda g_Y^2 - 9\lambda g_2^2 + 5\kappa^2 + \frac{5}{2}\kappa'^2 \right),$$

$$\frac{d\lambda_X}{dt} = \frac{1}{16\pi^2} \left(+36\lambda_X^2 + 48\lambda_X\lambda'_X + 720\lambda_X\lambda_X'' + 1152\lambda_X'\lambda_X'' + 3168\lambda_X''^2 + 6Y_X^4g_Y^4 - 12Y_X^2\lambda_Xg_Y^2 - 72\lambda_Xg_2^2 + 2\kappa^2 \right),$$

$$\frac{d\lambda'_X}{dt} = \frac{1}{16\pi^2} \left(+24\lambda_X\lambda'_X + 84\lambda_X'^2 + 408\lambda_X'\lambda_X'' - 84\lambda_X''^2 + 3g_2^4 + 12Y_X^2g_Y^2g_2^2 - 12Y_X^2\lambda'_Xg_Y^2 - 72\lambda'_Xg_2^2 + \frac{1}{2}\kappa'^2 \right),$$

$$\frac{d\lambda_X''}{dt} = \frac{1}{16\pi^2} \left(+8\lambda_X'^2 + 24\lambda_X\lambda_X'' - 32\lambda_X'\lambda_X'' + 368\lambda_X''^2 + 6g_2^4 - 12Y_X^2\lambda_X''g_Y^2 - 72\lambda_X''g_2^2 \right),$$

$$\begin{aligned} \frac{d\kappa}{dt} = & \frac{1}{16\pi^2} \left(+3Y_X^2 g_Y^4 + 18g_2^4 \right. \\ & + 12\kappa\lambda + 24\kappa\lambda_X + 24\kappa\lambda'_X + 360\kappa\lambda''_X + 6y_t^2\kappa \\ & \left. - \frac{3}{2}\kappa g_Y^2 - 6Y_X^2\kappa g_Y^2 - \frac{81}{2}\kappa g_2^2 + 4\kappa^2 + 6\kappa'^2 \right), \end{aligned}$$

$$\begin{aligned} \frac{d\kappa'}{dt} = & \frac{1}{16\pi^2} \left(+12Y_X g_Y^2 g_2^2 \right. \\ & + 4\kappa'\lambda + 4\kappa'\lambda_X + 60\kappa'\lambda'_X + 60\kappa'\lambda''_X + 6y_t^2\kappa' \\ & \left. - \frac{3}{2}\kappa' g_Y^2 - 6Y_X^2\kappa' g_Y^2 - \frac{81}{2}\kappa' g_2^2 + 8\kappa\kappa' \right). \quad (A.3) \end{aligned}$$

• SM with a sextet scalar field

$$\begin{aligned} \frac{d\lambda}{dt} = & \frac{1}{16\pi^2} \left(+24\lambda^2 - 6y_t^4 + \frac{3}{8}g_Y^4 + \frac{9}{8}g_2^4 \right. \\ & + 12\lambda y_t^2 - 3\lambda g_Y^2 - 9\lambda g_2^2 + \frac{3}{4}g_Y^2 g_2^2 \\ & \left. + 6\kappa^2 + \frac{35}{8}\kappa'^2 + \frac{28}{5}|\lambda_{\Phi^{\dagger 2}X^2}|^2 \right), \\ \frac{d\lambda_X}{dt} = & \frac{1}{16\pi^2} \left(+40\lambda_X^2 + 70\lambda_X\lambda'_X + \frac{3535}{2}\lambda_X\lambda''_X \right. \\ & + \frac{8525}{4}\lambda'_X\lambda''_X + \frac{271975}{16}\lambda_X''^2 + 6Y_X^4 g_Y^4 \\ & \left. - 12Y_X^2\lambda_X g_Y^2 - 105\lambda_X g_2^2 + 2\kappa^2 - \frac{11}{8}|\lambda_{\Phi^{\dagger 2}X^2}|^2 \right), \\ \frac{d\lambda'_X}{dt} = & \frac{1}{16\pi^2} \left(+24\lambda_X\lambda'_X + 136\lambda_X'^2 + \frac{2115}{2}\lambda'_X\lambda''_X \right. \\ & - \frac{265}{2}\lambda_X''^2 + 3g_2^4 + 12Y_X^2 g_Y^2 g_2^2 \\ & \left. - 12Y_X^2\lambda'_X g_Y^2 - 105\lambda'_X g_2^2 + \frac{1}{2}\kappa'^2 - \frac{7}{25}|\lambda_{\Phi^{\dagger 2}X^2}|^2 \right), \\ \frac{d\lambda''_X}{dt} = & \frac{1}{16\pi^2} \left(+8\lambda_X'^2 + 24\lambda_X\lambda''_X + 2\lambda_X'\lambda''_X + \frac{1715}{2}\lambda_X''^2 \right. \\ & \left. + 6g_2^4 - 105\lambda_X'' g_2^2 - 12Y_X^2\lambda_X'' g_Y^2 + \frac{2}{25}|\lambda_{\Phi^{\dagger 2}X^2}|^2 \right), \\ \frac{d\kappa}{dt} = & \frac{1}{16\pi^2} \left(+3Y_X^2 g_Y^4 + \frac{105}{4}g_2^4 \right. \\ & + 12\kappa\lambda + 28\kappa\lambda_X + 35\kappa\lambda'_X + \frac{3535}{4}\kappa\lambda''_X \\ & + 6y_t^2\kappa - 57\kappa g_2^2 - \frac{3}{2}\kappa g_Y^2 - 6Y_X^2\kappa g_Y^2 \\ & + 4\kappa^2 + \frac{35}{4}\kappa'^2 + \frac{56}{5}|\lambda_{\Phi^{\dagger 2}X^2}|^2 \right), \\ \frac{d\kappa'}{dt} = & \frac{1}{16\pi^2} \left(+12Y_X g_Y^2 g_2^2 \right. \\ & + 4\kappa'\lambda + 4\kappa'\lambda_X + 101\kappa'\lambda'_X + \frac{697}{4}\kappa'\lambda''_X + 6y_t^2\kappa' \\ & - \frac{3}{2}\kappa' g_Y^2 - 6Y_X^2\kappa' g_Y^2 - 57\kappa' g_2^2 \\ & \left. + 8\kappa\kappa' + \frac{64}{25}|\lambda_{\Phi^{\dagger 2}X^2}|^2 \right), \end{aligned}$$

$$\begin{aligned} \frac{d\lambda_{\Phi^{\dagger 2}X^2}}{dt} = & \frac{\lambda_{\Phi^{\dagger 2}X^2}}{16\pi^2} \left(+4\lambda + 4\lambda_X - 31\lambda'_X + \frac{961}{4}\lambda''_X \right. \\ & \left. + 6y_t^2 - \frac{3}{2}g_Y^2 - 6Y_X^2 g_Y^2 - 57g_2^2 + 4\kappa' + 8\kappa \right). \quad (A.4) \end{aligned}$$

• SM with a real septet scalar field

$$\begin{aligned} \frac{d\lambda}{dt} = & \frac{1}{16\pi^2} \left(+24\lambda^2 - 6y_t^4 + \frac{9}{8}g_2^4 + \frac{3}{8}g_Y^4 \right. \\ & + 12\lambda y_t^2 - 3\lambda g_Y^2 - 9\lambda g_2^2 + \frac{3}{4}g_Y^2 g_2^2 + \frac{7}{2}\kappa^2 \Big), \\ \frac{d\lambda_X}{dt} = & \frac{1}{16\pi^2} \left(+30\lambda_X^2 + 2448\lambda_X\lambda'_X + 51840\lambda_X''^2 \right. \\ & \left. - 144\lambda_X g_2^2 + 2\kappa^2 \right), \\ \frac{d\lambda''_X}{dt} = & \frac{1}{16\pi^2} \left(+6g_2^4 + 1530\lambda_X''^2 + 24\lambda_X\lambda''_X - 144\lambda_X' g_2^2 \right), \\ \frac{d\kappa}{dt} = & \frac{1}{16\pi^2} \left(+36g_2^4 + 12\kappa\lambda + 18\kappa\lambda_X + 6y_t^2\kappa \right. \\ & \left. - \frac{3}{2}\kappa g_Y^2 - \frac{153}{2}\kappa g_2^2 + 4\kappa^2 \right). \quad (A.5) \end{aligned}$$

• SM with a complex septet scalar field

$$\begin{aligned} \frac{d\lambda}{dt} = & \frac{1}{16\pi^2} \left(+24\lambda^2 - 6y_t^4 + \frac{3}{8}g_Y^4 + \frac{9}{8}g_2^4 \right. \\ & + 12y_t^2\lambda - 9\lambda g_2^2 - 3\lambda g_Y^2 + \frac{3}{4}g_2^2 g_Y^2 + 7\kappa^2 + 7\kappa'^2 \Big), \\ \frac{d\lambda_X}{dt} = & \frac{1}{16\pi^2} \left(+44\lambda_X^2 + 96\lambda_X\lambda'_X + 3744\lambda_X\lambda''_X \right. \\ & + 77184\lambda_X''^2 + 9216\lambda_X\lambda_X''' + 31104\lambda_X'\lambda_X'' \\ & + 90432\lambda_X''\lambda_X''' + 2859264\lambda_X''^2 + 6Y_X^4 g_Y^4 \\ & \left. - 12Y_X^2\lambda_X g_Y^2 - 144\lambda_X g_2^2 + 2\kappa^2 \right), \\ \frac{d\lambda'_X}{dt} = & \frac{1}{16\pi^2} \left(+24\lambda_X\lambda'_X + 204\lambda_X'^2 + 768\lambda_X'\lambda_X'' \right. \\ & + 7008\lambda_X''^2 + 44616\lambda_X'\lambda_X''' - 73824\lambda_X''\lambda_X''' \\ & + 2408676\lambda_X''^2 + 3g_2^4 - 12Y_X^2\lambda_X' g_Y^2 - 144\lambda_X' g_2^2 \\ & \left. + 12Y_X^2 g_2^2 g_Y^2 + \frac{1}{2}\kappa'^2 \right), \\ \frac{d\lambda''_X}{dt} = & \frac{1}{16\pi^2} \left(+8\lambda_X'^2 + 24\lambda_X\lambda''_X + 96\lambda_X'\lambda_X'' \right. \\ & + 1588\lambda_X''^2 + 912\lambda_X'\lambda_X''' + 5312\lambda_X''\lambda_X''' + 22696\lambda_X''^2 \\ & + 6g_2^4 - 144\lambda_X'' g_2^2 - 12Y_X^2\lambda_X'' g_Y^2 \Big), \\ \frac{d\lambda_X'''}{dt} = & \frac{1}{16\pi^2} \left(+16\lambda_X'\lambda_X'' - 80\lambda_X''^2 + 24\lambda_X\lambda_X''' \right. \\ & - 128\lambda_X'\lambda_X''' + 3056\lambda_X''\lambda_X''' - 12200\lambda_X''^2 \\ & \left. - 12Y_X^2\lambda_X'' g_Y^2 - 144\lambda_X'' g_2^2 \right), \\ \frac{d\kappa}{dt} = & \frac{1}{16\pi^2} \left(+3Y_X^2 g_Y^4 + 36g_2^4 + 12\kappa\lambda + 32\kappa\lambda_X \right. \\ & + 48\kappa\lambda'_X + 1872\kappa\lambda''_X + 4608\kappa\lambda_X''' + 6y_t^2\kappa \\ & \left. - \frac{3}{2}\kappa g_Y^2 - 6Y_X^2\kappa g_Y^2 - \frac{153}{2}\kappa g_2^2 + 4\kappa^2 + 12\kappa'^2 \right), \end{aligned}$$

$$\begin{aligned} \frac{d\kappa'}{dt} = & \frac{1}{16\pi^2} \left(+12Y_X g_Y^2 g_Y^2 + 4\kappa'\lambda + 4\kappa'\lambda_X \right. \\ & + 156\kappa'\lambda_X' + 384\kappa'\lambda_X'' + 13236\kappa'\lambda_X''' + 6y_t^2 \kappa' \\ & \left. - \frac{3}{2}\kappa' g_Y^2 - 6Y_X^2 \kappa' g_Y^2 - \frac{153}{2}\kappa' g_Y^2 + 8\kappa\kappa' \right). \quad (A.6) \end{aligned}$$

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